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NEW FUEL CELL ARCHITECTURES AND THE ROLE OF ELECTROKINETIC FLOWS IN ITS PERFORMANCE

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New Fuel Cell Architectures and the Role of Electrokinetic Flows in its Performance

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Abstract--This papers discusses the role of: transport phenomena, electrokinetic flows, gas-liquid interface area, convective mix, and the electrical boundary layer (EBL) in fuel cell performance. The electrokinetic phenomena analysis is based in a new model of electrocapillary flow, and its preliminary experimental tests. With this model, oxygen transport and convective ionic current in high specific area electrodes are analyzed. The paper also shows the use of the Critical Chain Method (CCM) in the appraisal of feasibility, costs and risks of innovative fuel cell set ups.

Index Terms-- Fuel cells, Fluid flow, Boundary Layers, Electrochemical devices, Electrochemical processes;

I. NOMENCLATURE

	0
A	- surface area
$[a_i]$	 concentration of dissolved ion i
a∞	- bulk ion concentration
\mathbf{B}_{i}	- influence coefficient
b _k	- constants in Grahame formula
С	- electrical capacity
c	- capacity per area
c _m	- capacity per area of dEBL
c_2	- capacitance per area of nEBL
EBL	- electric boundary layer
Eo	- Surface parallel electrical field
EKF	- electrokinetic flow
EOF	- electroosmotic flow
dEBL	- diffuse electric boundary layer
D_i	- diffusion coefficient
$\mathbf{E}_{\mathbf{x}}$	- Surface normal electrical field
F	- Faraday constant
h	- slip factor
Κ	- Inverse of Debyer length
nEBL	- non-diffuse electric boundary layer
Р	- Pressure
P	- Pressure far from surface
PB	- Poisson-Boltzmann
PEM	- Proton Exchange Membrane
R	- Universal gas constant
U	- electrical potential
U _w	- wall electrical potential
u	- fluid velocity

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u _m	- speed profile component
u _p	- electroosmotic speed component
Vp	- electrocapillary speed component
VOHP	- nEBL slip speed
х	- normal coordinate
z	- number of charges of an ion
α	- accommodation coefficient
δ	- conversion factor
ε	- media dielectric constant
$\phi_{\rm w}$	- fluid potential less U_w
γe	- volume charge density
φ	- non-dimensional potential
μ	- dynamic viscosity
ρ	- ionic resistivity
σ	- surface tension
σ_{0}	- pure liquid surface tension
ω	- fluid density
ζ	- potential inside EBL

II. INTRODUCTION

ach low temperature fuel cell has different transport phenomena restrictions. Standard plane gas-liquid interface PEM flooded models show that its peak power O2 entry by diffusion is limited to a few nanometers [1]. The use of test probes inside such PEM electrodes proved that only a tiny part of the catalyst layer is truly active [13]. PEM set ups with catalyst in its diffuser carbon filaments have no vital O2 transport limits, but strong ionic current drain. The Adams-Watson-Bacon fuel cell had, beside the problem of low O₂ solubility in alkaline liquids [2], great ionic current restrictions due to its electrodes low electrolyte volume fraction, air channels and particle scales tortuosities [11]. Ledoux et al. [3] used carbon fiber with covers of entangled carbon nanotubes as fuel cell electrodes. Bambace et al. [6,15,16] used tubular setups. MEMS based and micro-fluidic fuel cells and other non standard setups are also in the literature [5,6,12]. Such new setups have many Degrees of Freedom (DOF) that affect their operation, and we need to understand all effects to evaluate the potential advantages of such systems. In electroosmotic flow (EOF) chemical reactions control charge levels in the EBL as dissociation of surface silanol or PEM membranes sulfonic groups. Capacitive effects prevail in electrocapillary systems (metal-bath interfaces). Top speeds in 1-10 µm metal ducts can't be found with a simple capillary and pressure drop forces balances: the flow itself reduces its force drivers, as in thermocapillary flow. The Knudsen number (Kn) based on EBL thickness is 0.03-0.15, so Kn might be in the slip flow range. Surface diffusion and the mobility of adsorbed ions are known to reduce the roughness of electroplated parts. So, the non-slip condition is not valid, and experiments are needed to obtain new flow models. Convective mixing may appear in some cases due either to Bénard cells or tortuous flows induced by random nanometer features.

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Convection helps ionic current transport only if the electrolyte is not neutral, i.e., inside the EBL, which is thin to reduce electrical potential energy. Onsager's principle [7] implies that an electrokinetic flow (EKF) only appears if it reduces the system entropy generation. So the power required to keep EKFs comes from a fall in another kind of dissipation. There are difficulties in achieving reliable and repetitive building processes of new setups. So, feasibility and risks of innovative projects was studied with a version of the Critical Chain Method that can rank research options, not only hinting the need of doing any option, but saying when to start it, based on the outputs of already scheduled activities.

III. MODELS DESCRIPTION

Static electrocapillary was described in the 19th century by Lippmann, who noticed that mercury surface tension had a parabolic profile with a maximum for uncharged liquid. For molecules with dipoles in the electrolyte, the surface tension σ is given by:

$$\sigma = \sigma_{o} - 0.5 (C/A) \Delta U^{2} + \sum_{i \in B_{i}} \ln(1 + b_{k} [a_{i}])$$

$$(1)$$

where σ_0 is the uncharged pure liquid surface tension, C the electrical capacity, A the area, U the electrical potential, $B_i \in R$, an influence coefficient, $[a_i]$ the concentration of dissolved ion i, and b_k are constants. The energy is the $d\sigma/dx$ integral over a surface, so capillary force in any channel of fixed perimeter p, with tip σ values being σ_1 and σ_2 and $p(\sigma_1-\sigma_2)$. For a tube of diameter d, length L, and fluid dynamic viscosity μ , the force balance is given by the flow speed of $(\sigma_1-\sigma_2) L^{-1} \mu^{-1} d/8$ if standard speed profiles are used. Large tubes don't show such speeds no matter what causes the changes of σ . An EBL of thickness x_0 has a non-diffuse part of thickness where dipoles forces prevail. It also has a diffuse sublayer (dEBL) that obeys an one-dimensional Gauss equation, as Poisson-Boltzmann (PB) so for $\phi = U E/(R T)$, and q=1:

$$\nabla^2 \varphi = \varepsilon^{-1} \gamma_e = -K^2 \Sigma n_i n_o^{-1} \exp(-q z_i \varphi)$$
⁽²⁾

where the potential reference is set far away from the surface, a place with ion concentration a_{∞} , F is the Faraday constant, R the universal gas constant, T the absolute temperature, K is $[2z_i^2F^2a_{\infty}(y)]^{1/2}[\epsilon R T]^{-1/2}$ for y a coordinate parallel to the surface, z_i the number of charges of ions, γ_e the volume charge density, ϵ the media dielectric constant. Replacing Boltzmann distribution other equations arise [8-9], generally written in similar form with other q and K, at least as an approximation. The solution of equation (2) is often written for constant $z_i=z$ in implicit form:

$$tanh(0.25 \varphi) = tanh(0.25 \varphi_2(y)) exp(-K (x-x_2))$$
 (3)

where $\varphi_2(y)$ is the dimensionless potential at x_2 and x the distance from the surface. Equation (3) may be rewritten, with help of $\delta = z F/(R T)$ and $\zeta(y) = \varphi_2(y) \delta^{-1}$ in explicit form as:

$$G(x,y) = \tanh(0.25 \ \delta \ \zeta(y)) \exp(-K \ (x-x_2)) \tag{4.a}$$

$$U(x,y) = 2 \delta^{-1} \ln\{1 + G(x,y)\} - 2 \delta^{-1} \ln\{1 - G(x,y)\}$$
(4.b)

Unit area capacities $c \equiv (C/A)$ may be split in diffuse and non-diffuse parts, $c_m = (x_o \cdot x_2)/\epsilon$ and $c_2 = x_2/\epsilon$ PB fails to model cathode and anode c differences. In the EBL there is an electrical field \mathbf{E}_o in the direction $\hat{\mathbf{e}}_y$ and a much stronger field \mathbf{E}_x in the direction $\hat{\mathbf{e}}_x$ normal to the metal surface, where charge density v is $[8 \, RT\epsilon \, a_{\infty}]^{1/2} \sinh(0.5 \, \delta \zeta)$ in PB. U and ionic resistivity ρ obeys the Laplace Equation, thus for E the electrical field:

$$\mathbf{c}^{-1} = \{ \varepsilon^{-1} [\cosh(\delta \zeta(\mathbf{y})/2)]^{-1} \mathbf{K}^{-1} + \varepsilon^{-1} \mathbf{x}_2 \} = \mathbf{c}_2^{-1} + \mathbf{c}_m^{-1}$$
(5)

$$\mathbf{E} = -\mathbf{V}\mathbf{U} \tag{6}$$

$$\nabla \rho \ \nabla U = 0 \tag{7}$$

As the layer is thin $\gamma_{e} \sim \epsilon d^2 U/dx^2$. Continuity, Navier Stokes and species concentration equations are:

$$\nabla_{\mathbf{v}} \mathbf{u} = \nabla_{\mathbf{v}} (\mathbf{u}_{\mathbf{x}} \, \hat{\mathbf{e}}_{\mathbf{x}} + \mathbf{u}_{\mathbf{y}} \, \hat{\mathbf{e}}_{\mathbf{y}}) = 0 \tag{8}$$

$$\omega \mathbf{u}_i \, \mathbf{d} \mathbf{u} / \mathbf{d} \mathbf{x}_i = -\nabla \mathbf{P} + \mu \nabla^2 \mathbf{u} + \gamma_e \mathbf{E} \tag{9}$$

$$-\mathbf{D}_{i} \nabla^{2} \mathbf{a}_{i} + \mathbf{m}_{i} \nabla^{\circ} (\mathbf{E} \mathbf{a}_{i}) + \nabla^{\circ} (\mathbf{u} \mathbf{a}_{i}) + \chi_{qi} - \chi_{qo} = 0$$
(10)

where ω is the fluid density, a_i species i concentration, χ_4 and χ_{Ψ} chemical conversion rates, \mathbf{m}_i ionic mobility, D_i diffusion coefficient. EBL speed component normal to a solid is negligible, and, 2 nm< x_0 <5 nm, to P_{ω} the far away pressure, pressure is $P(xy)=P_{\omega}(y)+\epsilon (dU/dx)^2/2$. Due to adsorbed ions surface diffusion and mobility of, and a possible slipping boundary condition, a tube fluid velocity \mathbf{u} may be split in 4 wall parallel parts: a nEBL slip v_{OFP} at $x_2, u_p = (\zeta - U)\epsilon \mathbf{E}_0 \mu^{-1}, u_m = U_m(r^2/R^2)$, and v_p ; and \mathbf{s} normal to wall.G=0 out of EBL,dU/dx=0 if U=0, so:

$$[\varepsilon d^{2}U/dx^{2}E_{o}-\mu d^{2}u_{p}/dx^{2}]+$$
+
$$[dP/dy+4\mu U_{m}R^{-2}+\Delta P_{a}+\omega u du/dy +\omega s du/dx] -$$
-
$$[\mu d^{2}v_{p}/dx^{2}-\varepsilon dU/dx d^{2}U/dxdy-\varepsilon dU/dy d^{2}U/dx^{2}]$$
(11)

Setting the 1st line of equation (11) to zero, one gets the classical EOF equation, solved by the u_p , the 2^{nt} line governs flow out of the EBL, and is solved by u_m if s=0. The 3st bracket refers to the electrocapillary extra component v_p , function of electrical variables only. For x-x₂<x_o-x₂ and $L=2\mu^{-1}\epsilon\delta^{-2}a_{o}^{-4}da_{o}/dy$:

$$dv_p/dx = \mu^{-1} \epsilon dU/dy dU/dx$$
(12)

$$Y = -2 \mu^{-1} \epsilon \delta^{-1} \operatorname{cosech}(\delta \zeta/2) (d\zeta/dy)$$
(13)

$$v_p = [1-G(x,y)^2]^{-1} \{Y-L K(x-x_2)\} +$$

$$+ L \ln[(1+G(x,y))/(1-G(x,y))] \tanh(\delta \zeta/4)^{-1}$$
(14)

See that a_i and a_{∞} vary in finite gaps due to dissociation changes that overcome migration of non-reacting ions. For 2D cases with a fixed wall potential U_w , $\phi_w \equiv U_w$ -U is the total EBL cross potential, at $\phi_w = 0$, $d\phi_w/dy = (x_2 \text{ K}+1) d\zeta/dy$, so for linear U profile outside a PB EBL,with K and capacitance definitions, it is possible to write:

$$d\phi_w/dy = -0.5 (U_w - \zeta) a_{\infty}^{-1} da_{\infty}/dy$$
$$-0.5 \delta U_w [tanh(\delta \zeta/2)]^{-1} d\zeta/dy + 0.5 \delta \zeta tanh(\delta \zeta/2) d\zeta/dy$$
(15)

The integral of $\gamma_e(u_p+u_m+v_p+v_{OHP})$ in the $\hat{\mathbf{e}}_x \hat{\mathbf{e}}_z$ planes is the flow current. As dEBL is thicker for low ϕ_w the peak u_p+v_p is in low ϕ_w zone, where flow current is minimal. Ignoring EBL area in mass balance $U_m(y) = \ddot{U}_m \cdot \sum \{v_{p\infty}(y) + u_{p\infty}(y) + v_{OFP\infty}(y)\}/2$, and integrating 2^{rd} line of equation (11) for given tube end pressures, one gets \ddot{U}_m . For tubes of $10^5 m$, U profile is near linear, to 10⁸ m systems, flow current is high and this ionic potential drop outside EBL is nonlinear, so iterative schemes and $\hat{\mathbf{e}}_{v}$ interpolation are used to find dU/dy. Equation (7) is solved only inside EBL, and out of it real profile approximated with straight line segments, as for each grid element current is assumed constant. VOHP minimizes the sum of plastic all dissipations, including the nEBL plastic one if ϕ_w is high. For EBL thickness with small face grid scales, it is wise to supply the total EBL slip to the CFD model as a function of local ϕ_w . In general electrode potentials, total current and chemical kinetic models the are known. EKF test set ups have one channel and connected metal surfaces in short circuit, so the ϕ_w profiles minimize capacitive energy. Reactions may occur in insulated metallic parts, as tips work either as a secondary cathode or anode if

 $d\phi_w/dy$ is high. For **n** molecules by unit of volume with mean cross section diameter ξ , the mean free path λ is $(2^{0.5}\pi n\xi)^{-1}$. Reference [10] hints a slip $(2-\alpha)\mu^{-1}\alpha^{-1}\lambda\tau_x+.75(\gamma-1)\gamma^{-1}\mu\omega^{-1}T^{-1}dT/dx$. Its dT/dx part (ψ) is the thermocapillary slip, τ_x is the hydrodynamic stress $\mu du_m/dx$, α and γ accommodation coefficients. For high ϕ_w , x_0 - x_2 and ζ vanish. So with this slip as example we took $v_{CHP} = (2-\alpha)\mu^{-1}\alpha^{-1}[\mu\lambda d\nu dx + h\epsilon c\phi_w E_o]$, with h=0 for $\phi_w < \phi_c$, with nEBL yield strength reached at ϕ_c . EBL forces may bias the molecules shifts due to thermal shocks. Thus, it may happen that h≠0always, v_p and u_p may be high in some points, but cathodic and anodic values are opposed to each other, and don't explain EKF flow for asymmetric ϕ_w fields. Equation (14) is not true at $\phi_w=0$, where a wall outward wall flow can't be ignored, $s \neq 0$. v_p and v_{CHP} main driver is pressure change, there is an electrostatic pressure due to charge unbalance in EBL. Without dEBL, it is easy to see that the nEBL pressure gradient times its thickness is exactly the electrocapillary force. The total pressure force in the EBL is always equal to the electrocapillary force, although it is not so easy to prove. Multiplying ψ by $u\omega d\mu^{-1}$, it's easy to see that ψ is a speed.

IV. EXPERIMENTAL VERIFICATION

The theoretical model and an $\alpha=1$ hypothesis were tested with the setup showed in Figure 1. Electroplating tapes were placed to avoid unpredicted passages due to distortion of PVC parts. The potential difference cross the 2 pools gives passages ionic current. Applied current was measured directly. Both currents are equal if applied potential is less than twice the electrolysis one, if not metal channel metallic annexes work either as a secondary anode or cathode. Ionic resistivity ρ was found with fixed current, i, and 2 smooth gold plates of area S in a rectangular box. Moving a plate, only the current path L changes, so steady state applied potential change is i $\rho\Delta L/S$. The rest of applied potential is due to activation and boundary layer polarizations (see figure 2). All parts were jointly electroplated, with distinct cathode current controls. Due to some degree of porosity and corrosion, it was not possible to do alkaline tests, or to replace the tape barrier by more regular one. Acid tests were filmed with applied potential at main electrodes from 1.4 to 3.2 V, in 0.2 V steps. Bubbles hit one another in the channel and join themselves in bigger ones, as well as arrive at the fluid upper surface, or hit the walls. In 63 s films, 3 to 5 bubbles that pass directly through the entire channel were found. Events were observed frame by frame with the VirtualDub software, to obtain the average bubble speeds in the channel to each applied potential. The selected frames were exported to JPEG format with this program. Figure 3 shows 5 frames took from one of the films that shows a bubble passing the channel. Average speeds are proportional to average surface tension force. As the capacitive energy of channel and surrounding metal parts is minimal, and as anodic zone has a limit capacitance of 0.2 F/m^2 , and cathodic zone a 0.7 F/m^2 one, channel tip potentials were found to be at a rate of 0.37 in modulus. Total channel potential were measured with Pt 36 AWG wires. With this we found an **h** value of 57 ± 15 . Errors come from oxidation, bad cleaning, corrosion rate and errors in electrical measurements, as nEBL thickness of clean surfaces are less than 0.3 nm. Cleaning is decisive to EKF intensity level.

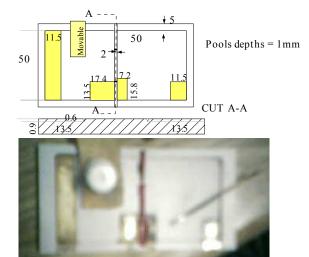


Fig 1: Flow test Arrangement.

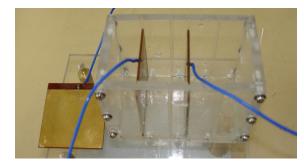


Figure 2: Ionic conductivity test set up.

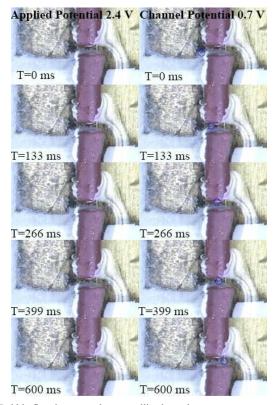


Fig. 3 Bubble flow in a very clean metallic channel

TABLE 1: SUMMARY OF MAIN TEST DATA FOR H_2SO_4 Solution.

Total Potential (V)	Channel Potential (mV)	Current (mA)
1.40	166	1.15 <u>+</u> 0.15
1.60	371	3.5 <u>+</u> 0.3
1.80	563	12.5 <u>+</u> 0.3
2.00	754	18.4 <u>+</u> 0.3
2.20	959	22.3 <u>+</u> 0.3
2.40	1.21	29.3 <u>+</u> 0.2
2.60	1.32	36.7 <u>+</u> 0.3
2.80	1.57	42.3 <u>+</u> 0.1
3.00	1.76	50.0 <u>+</u> 0.1
3.20	1.87	57.8 <u>+</u> 0.3

V. NUMERICAL RESULTS

A common Pt catalyst membrane painted PEM was studied with the flooded model to generate a baseline solution. It has a very weak EKF, as contact resistances and ionomer insulation limits ϕ_w to small values. Later, we studied the fuel cells of Figures 8 to 11: the Bacon cell, the Ledoux carbon nm-tube and fiber plies PEM, and 2 µ-tube cells. In these 4 cells, the gas goes from a high pressure to a low pressure zone, passing through the pores or through a 3 mm thick non-woven arrangement of tubes or coated plies, separated by large pore foams with end thickness of 0.2 and 0.8 mm. Preliminary theoretical results for flow current in non-reacting Ni u-tubes with nm-foam walls, with the acid h value instead of alkaline are in table 2. Figures 12-16 give theoretical results for: output potential, efficiency and power for each current without^a or with^b EKF. The exchange current for the Ni tubular cell, Figure 14, was 3.10⁵ A/m², 35% in volume of carded tubes of 5 and 7 μ m radius, 100 nm foam texture with 85% of void. PtO₂ reduction exchange current is 0.07 A/m². For * marked case, dEBL gives 28% of flow current. CFD analysis predicted 6 to 9 Bénard cells for the arrangements of Figures 9-11, with fibers normal to ionic insulation. The strong flow out of the wall near the ionic zero is an inducer of such cells with lengths ranging from 30 to 150 µm. It shall be seen that even without EKF, performances of cylindrical gas-liquid interface cells are higher.

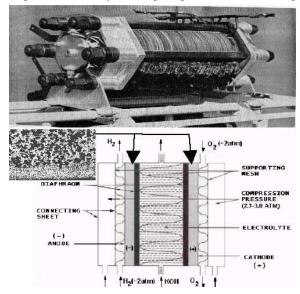


Figure 8: Adams-Bacon-Watson Alkaline Fuel Cell.

CFD analysis predicted 6 to 9 Bénard cells in figures 9-11 set ups, for fibers normal to ionic insulation. The strong flow outward the wall near the ionic zero induces of such cells with lengths ranging from 30 to 150 mm.

TABLE 2: FLOW CURRENT DATA FOR NI M-TUBES WITH 6 M KOH.

Applied Potential	Total Lenght	Null ø _w Point	Peak Flow Current	Meam Bulk Current
(V)	(mm)	(mm)	(µA)	(µA)
0.1	0.80	0.10	0.0487	0.0188
0.2	0.80	0.10	0.383	0.0377
0.3	0.80	0.10	1.29	0.0565
0.4	0.80	0.10	3.07	0.0754
0.1	0.80	0.40	0.0335*	0.0188
0.2	0.80	0.40	0.224	0.0377
0.3	0.80	0.40	0.755	0.0565
0.4	0.80	0.40	1.79	0.0754

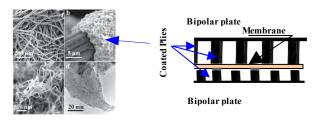


Fig. 9: Ledoux PEM Fuel Cell.

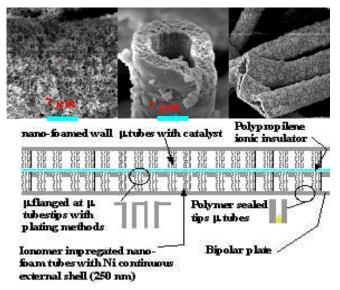


Fig. 10: Nickel nanofoam wall µ-tubes AFC.

The results for cells in Figures 9 and 11depend strongly on the density of the nanotube electric connections and related resistance. It will be seen that even without EKF, performances of cylindrical gas-liquid interface cells are better. Flow current is the major cause for better results of all cells of figures 8 to 11 when they are clean and EKF is present. Tubular cells have better ionic transport, and Bénard cells that slightly enriches their tube cores with O_2 , creating a 2nd diffusion path and about 2% more O_2 flux. The foarned tubular cell has random foam surfaces, forcing zig-zag flow, enhancing O_2 transport in 10%, approximately; its current is smaller than Ledoux and tubular acid cells, due to the low O_2 solubility in alkali, but it is a cheaper device. The expenses with noble metal catalyst are the major part of an acid fuel cell cost. Acid corrosion resistance of the cell presented in Figure 10 covered with NiCrN as Los Alamos metal bipolar plates is a doubt.

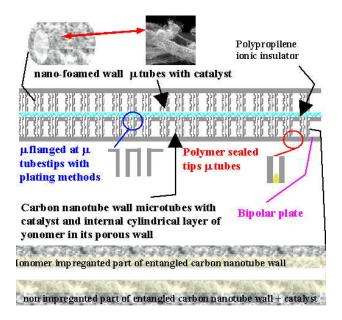


Fig. 11: Carbon nanotube wool wall m-tubes fuel cell.

Cell performance curves were based on chemical kinetic models and on energy conservation, so output voltages are the open cell less the integral of all losses, divided by the currents. O₂ and H₂ local concentration, catalyst data and diffusion flux limits were used to calculate local reaction rates and output currents. Total tubes area is 4 times their actual volume over their diameter. The cell in Figure 11 has more breathing area and available section to ionic conduction than the Ledoux cell and no membrane so its results are better, despite of no flow current in the anode. The cell of Reference [11] has porous carbon tubes with gas inside, catalyst and Nafion outer coats, and tubes are in an electrolyte pool, having power over 8 W/cm³, despite of its tubes separation. O₂ overpotential η is RT $\ln\{i^2 [a_0]^2 4FD_{eff} i_0 \kappa\}/(4F\alpha)$, where j is the current per unit area of the gas-liquid interface, $[a_0]$ the reference O₂ chemical activity, F the faraday constant, Def the effective diffusion coefficient of porous media, R universal gas constant, T the absolute temperature, io the exchange current, κ the area density, α the transfer coefficient. Common PEM cells main loss source is η , and η falls drastically in large area systems of small j. PEM electrodes have also large electrical resistance due to contact resistance between carbon black particles and meaningful membrane ionic resistance. Cylindrical geometries replace exponential solution by an alike eigenvalue I₀ Bessel function. Uniform reaction rate tubes have an effective resistance proportional to 1/3 of its length, and real length loss is smaller as reaction rates drops as the distance from ionic bridge rises.

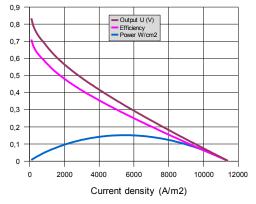


Fig. 12: Simple coated membrane cell performance.

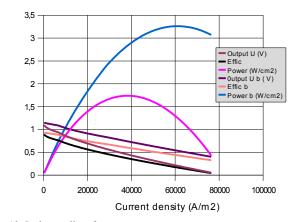


Fig. 13: Ledoux cell performance

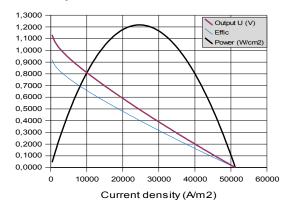


Fig. 14: Adams-Bacon-Watson AFC performance.

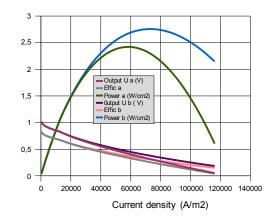


Fig. 15: Nickel nanofoam tubular cell performance.

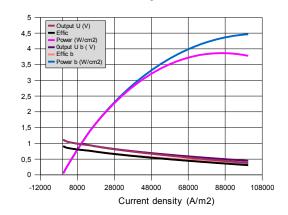


Fig. 16: Nanotube wool wall tubular cell performance.

VI. CRITICAL CHAIN METHOD IN PLANNING

The amount of research on fuel cell research is huge and most efforts are made without any coordination. This policy directs undue efforts to some fields, with many people doing nearly the same work, with lack of assets on other areas. This issue is less crucial in rich than in developing countries. All research projects hava risks. People may not reach desired goals, or a poor quality product may be developed. Research has lots to do with gambling, and good gamblers consider game trees. The delivery of very innovative products in military bids on time, assuring strategic advantages, requires more management expertise than the use of game trees or PERT schemes can provide. The Critical Chain Method (CCM) [14] may be improved with risk control tools employing parallel developments, if needed, judicious decisions about starting or not any extra task, and schedule [16], leading to CCM*. To control a risk, people must know that it exists, so brain storming shall be used to enlist possible troubles and to avoid failures in risk detection. Also the larger the number options studied, the greater the chance to work on a real good one. Fuel cell projects require choices of: geometry, catalysts, materials, gas-liquid interface control, ionic and electrical insulators, as well as fuel and oxidant sealing, separation and supply. Real task trees are too big to detail here, as single subsystems involve many choices. To explain CCM*, system level trees are given in figures 18-20, where rectangles, losangles and circles, match tasks, baffles and decisions listed in Tables 3-5. Failure Risk (FR) and Task Duration (TD) in weeks, are in the same tables. Comparing tables 3 and 5 for PEMFc systems it is seen that the main success factor is a good detailing of options and risks, which is crucial to innovation and market advantage. The risks of project failure for tubular cells and PEMFC are similar, and differences may be either associated with trees detailing or a real warning of more difficulties. But a tubular cell development needs much more effort and money, the cost of a bigger chance of innovations and a big market in crucial economic sectors. The improvement of techniques is difficult. For instance, to apply common 4-6 nm Pt particles of over any substrate is a low risk task, but there's a lot of risk of not applying 2 nm size Pt particles. Decision in CCM* gets the standard treatment given in CCM for external constraints.

TABLE 3: NO CHOICE TASKS CHAIN FOR MEMBRANE COATED PEM

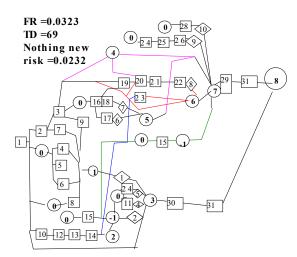
#	Task / Decision	FR	TD					
1	Making a new membrane	0.30	18					
2	Developing a system to apply smaller Pt grains	0.30	16					
3	Buy bipolar separator, diffuser and faulting parts	0	3					
4	Integrate the cell	0.02	1					
5	Test thecell	0	3					
#	Weeks# Weeks #							
1	2 2 2 13							
1	FR = 0.1082 3 4 5 2 TD = 29 Nothing new ris	sk =0.0)9					

Fig. 18: No choice PEM developing program

Initial studies showed that less active organometalic catalysts may be used due to better O_2 entry. Some of these catalysts stop making four O_2 electrons reduction at high polarization. A more active catalyst in H_2O_2 reduction may be used in the inner part of the tube wall to take advantage of convection. The lack of Pt free acid catalysts data is the main difficulty for the simulation of such designs. The wall total force equals electrical forces in steady state boundary layers, pressure field distributes and relieves the loads, but good quality control is needed to avoid Ni tubes fatigue.

TABLE 4: TUBULAR FUEL CELL TASKS CHAIN.

#		Tas	k / Dec	isio	n/doubt	*=>	ok	FR	TD
1	Task / Decision/doubt *=>ok Making Ni nanofoam wall µ-tubes mec. resistance						0.05	16	
2	Cr covering of task 1 µ-tubes cover uniformity						0.10	12	
$\frac{-}{3}$	Nitretation of task 2 μ -tubes composition-thermal treat.							0.01	3
4			•				wall possible	0.20	16
5	Applying						_	0.40	16
6		-	-				ansp dynam	0.10	16
$\frac{0}{7}$	Making N		-					0.10	8
8	Making L		· ·					0.05	8
9	Putting Pt		-					0.08	6
10	Nafion im					-		0.05	6
11	Buying po							0	1
12			-				ds porosity	0.05	8
13	_						obstructions	0.03	12
14	Put chemi	cal N	i ext part	imp	regnated]	Ni tub	e uniform	0.05	4
15					-		delamininate	0.30	18
16	Making ca	arbon	nanotub	e wa	ill µ-tubes	axial	coupling	0.07	22
17							/20 tubes)*	0.02	2
18	Organom	etalic	catalyst i	n tas	k 16 tube	s (8 op	ot) several	0.12	11
19	Nafion Im	pregr	nation of	task	16/3 tube	s volu	me controls	0.08	5
20	Sealing tas	sk 161	ubes bipo	olars	separator s	ide en	ds porosity	0.12	4
21	Making in	sulato	rside µ-f	lang	es in task 2	20tube	s obstructions	0.03	12
22	Applying	Pt to t	ask 21 tu	ibes	*			0.0	2
23	Applying	Ru in	ner coat i	tube	s task 12 c	or 13 ti	ubes several	0.03	1
24	Buying 18 µm Ni foil *							0	1
25	Cr over task 24 foil *							0	2
26	Nitretation task 23 foil with Los Alamos Patent allowed?							0	2
27	Put chosen catalyst in chosen acid cathode *(if cat ok) 0 2							2	
28	Carbon bipolar separator acquisition * 0 3						3		
29	Integration acid cell tools/cleaning 0.03 4								
30	Integration of alkaline cell tools/cleaning						0.03	4	
31	Validation							0	3
-1	Kill- no								
0	Start / n								
1	Alkaline				-				
2	Alkaline anode type (nafion in Ni wall or not)								
3	μ-flange / membrane or not alkaline mount								
4	Acid cathode support choice								
5	Acid cathode catalyst choice								
<u>6</u> 7	Acid anode type: NiCrN/C impreg. / membrane								
7	Acid mounting type Cell choice								
8 #	Cell cho Wæks	ice #	Weeks	#	Weeks	#	Weeks	#	Weeks
# 1	Weeks 8	# 2	16	# 3	14	# 4	1	# 5	1
$\frac{1}{6}$	° 23	2 7	27	3 8	14	9	1 2	10	3
<u> </u>	20	′	21	0	11	1	<u></u>	10	<i>.</i>



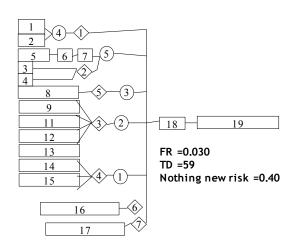


Figure 20: Good PEM Program Task Tree.

VII. CONCLUSIONS

The intensive development effort in PEM fuel cells lead some of its configurations performances very close to theoretical limits. Significant improvements in planar electrodes cells are very difficult. Jumping from gas liquid interface areas from a range of 1-10 times the nominal separators area to levels over 50, are still risky, but acceptable ones face performance increases, as shown with the Critical Chain Method. The raise in effective gas-liquid interface by unit of volume is the most important path for better performance, but effective increasing requires designs with good ionic current transport. Electrokinetic effects may have a key role in total ionic resistance, being able even to decrease them by 50%. O₂ transport improvements are mainly due to higher gas-liquid interface areas. Electrical Boundary Layer shall be better studied with other models, as the Poisson-Nerst-Planck, or models with ions size and relative positional electrostatic energy corrections. Ionic resistance is minimized with high ionic potentials and flow current near the ionic bridge boundary, where reaction rate is greater. High gas liquid interfaces dismiss the need of very effective catalysts. The joint use of better geometries and new catalysts will make fuel cells able to deliver more than 20 W/cm³ with efficiencies over 50% in about a decade, with great impacts on the energy market.

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Fig. 19: Tubular Fuel Cell Task Tree.

TABLE 5: GOOD TASKS CHAIN FOR MEMBRANE COATED PEM.

#									FR	TD
1	P ₂ O5 ⁻² /WO ₃ ⁻¹ composite membrane							0.30	18	
2	Silico-tungstic composite membrane								0.30	20
3	Perfluric ic	nomer co	vei	for con	nposi	te memb	rane		0.10	16
4	PEEK ion	omer cove	r f	for comp	osite	membra	ne		0.40	25
5	Anodized	carbon n	m-	tube scre	en				0.07	17
6	Filling with	n ionomer	tas	sk 5 scre	en				0.03	3
7	Activation	of task 6	pa	ırt expos	ed ca	rbon			0.03	3
8	Los Alamo	s diffuser v	vitl	n yarns ti	o hok	l more ior	nomer		0.15	22
9	Alloy catal	yst 1 for C	CH	$_4O/C_2H$	0_{6}				0.08	28
10	Alloy catal	yst 2 for	CH	I_4O/C_2H	I ₆ O				0.08	28
11	Alloy catal	yst 3 for	CH	I_4O/C_2H	I ₆ O				0.08	28
12	Alloy catal	yst 4 for	CH	I_4O/C_2H	I ₆ O				0.08	28
13									32	
14	Platted metal bipolar separator 2 0.30 32								32	
15	Platted metal bipolar separator 3 0.30 32								32	
16	Alcohol feeding 0.03 16									
17	O ₂ feeding 0.03 12									
18	Integration 0.03 2									
19	Validation 0 3							3		
-1	Kill- no ki	l task								
0	Start / no s	tart task								
1	Bipolar pla	ate								
2	Catalyst									
3	Conventional new diffuser									
4	Membrane base material									
5	Membran	e cover			_			_		
#	Weeks #	Weeks	#	Weeks	#	Weeks	#	Wee	ks	
1	10 2	25	3	18	4	20	5	22		
6	16 7	12	8							

Buying parts in the market reduces failure risks to a near null level, but these parts brings no innovation to the design. Carbon systems are intrinsically stronger but more expensive. Cleaning is the major predicted difficulty to take advantage from flow current, as if non-diffuse EBL is thicker than its 0.3 nm, its capacity decreases, and total charges also decreases.

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