## The sinusoid as the longitudinal profile in backward-wave oscillators of large cross sectional area

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#### Abstract

This work reports on design tools for BWOs operating in symmetric TM modes since these modes are able to perturb the axial velocity and electron density density on rectilinear beams confined by an external magnetic field in slowwave systems. Here we investigate whether a cylindrical guide with sinusoidally rippled wall can provide strong coupling between the guide surface waves and mildly relativistic (500-keV) electron beams in the 8-9 GHz frequency range for BWOs of large diameter ( $\mathrm{D} \sim 3 \lambda$ ). For this purpose, the characteristic equation of a sinusoidally corrugated structure is derived on the basis of the Rayleigh-Fourier method, whereby the field solution is represented by a single expansion of TM eigenmodes. From the dispersion diagrams thus obtained we infer the appropriate periodic length and ripple amplitude of the guiding structure that optimize the beam-wave interaction.


## DISPERSION RELATION OF THE CYLINDRICAL GUIDE WITH SINUSOIDALLY RIPPLED WALL

WAVEGUIDE PROFILE: $\quad R_{w}(z)=R_{0}\left(1+\varepsilon \cos \left(\frac{2 \pi}{d} z\right)\right)$
FLOQUET'S THEOREM :

$$
\begin{gathered}
E_{z}(r, z, t)=\sum_{n=-\infty}^{\infty} A_{n} J_{0}\left(k_{\perp n} r\right) \exp \left(i k_{z n} z-i \omega t\right) \\
E_{r}(r, z, t)=-i \sum_{n=-\infty}^{\infty} A_{n} \frac{k_{z n}}{k_{\perp n}} J_{1}\left(k_{\perp n} r\right) \exp \left(i k_{z n} z-i \omega t\right) \\
k_{z n}=k_{z 0}+\frac{2 \pi}{d} n, \quad 0 \leq k_{z 0} \leq \frac{2 \pi}{d} \\
k_{\perp n}^{2}=\frac{\omega^{2}}{c^{2}}-k_{z n}^{2}
\end{gathered}
$$

## Boundary Condition:

$$
\left[E_{z}(r, z)+E_{r}(r, z) \frac{d R_{w}(z)}{d z}\right]_{r=R_{w}}=0
$$

Homogeneous Equations:

$$
\begin{aligned}
& \sum_{n=-\infty}^{n=\infty} A_{n} C_{n m}\left(\omega, k_{z 0}, \varepsilon, R_{0}, d\right)=0 \\
& C_{n m}=\left[1+\frac{2 \pi}{d}(n-m)\right] \frac{k_{z n}}{k_{\perp n}^{2}} \int_{0}^{d} d z J_{0}\left(k_{\perp n} R_{w}(z)\right) \operatorname{Cos}\left((n-m) \frac{2 \pi}{d} z\right.
\end{aligned}
$$

Dispersion Equation:

$$
\operatorname{det}\left\|C_{n m}\right\|=0
$$



Fig. 1 BWO schematics and dispersion diagram showing the first passband


Fig. $2 \mathrm{TM}_{01}$-mode upper cutoff frequencies ( $\pi$-point) as function of the ripple amplitude $\varepsilon$ with the normalized period $\mathrm{d} / \mathrm{R}_{0}$ as parameter


Fig. $3 \quad \mathrm{TM}_{01}$-mode phase velocity (at the $\pi$-point ) as a function of the ripple amplitude $\varepsilon$ with the normalized period $\mathrm{d} / \mathrm{R}_{0}$ as parameter


Fig. 4 Sinusoidal $\left(R_{0}=4.5 \mathrm{~cm}, \varepsilon=1 / 15, R_{\text {min }}=4.2 \mathrm{~cm}, R_{\text {max }}=4.8 \mathrm{~cm}\right)$ and piecewise (semicircle radius 0.5 cm and rectangle height $h=0.1 \mathrm{~cm}$ ) profiles with periodic length $d=1.4 \mathrm{~cm}$


Fig. $5 \mathrm{TM}_{01}$-mode dispersion diagram corresponding to the sinusoidal profile in Fig. 4. Dashed and dotted lines are the light line and the $500-\mathrm{keV}$ beam Doppler line ( $\mathrm{v}_{\mathrm{z}}=0.86 \mathrm{c}$ )


Fig. 6 Radial dependence of the $\mathrm{TM}_{01}$-mode electric field on section $\mathrm{z}=0.0$ of sinusoidal and piecewise profiles shown in Fig. 4


Fig. $7 \quad \mathrm{TM}_{01}$-mode dispersion diagram for an oversized guide with longitudinal profile sinusoidally corrugated. The four curves are associated with distinct values of average radius $\mathrm{R}_{0}$ and ripple amplitude $\varepsilon$, but keeping in all cases $R_{\text {min }}=4.2 \mathrm{~cm}$ and $d=1.4 \mathrm{~cm}$. Dashed line is the light line.


Fig. 8 Radial dependence of the $\mathrm{TM}_{01}$-mode electric field (on section $z=0.0$ ) in sinusoidally rippled waveguides


Fig. 9
Calculated dispersion diagrams for three lower-order $\mathrm{TM}_{01}$-modes in the sinusoidally corrugated waveguide ( $R_{0}=4.5 \mathrm{~cm}, \mathrm{~d}=1.4 \mathrm{~cm}$, $\varepsilon=1 / 15$ ) shown in Fig. 4. Light and beam lines are indicated by dotted and dashed lines.

| h(cm) | WaveSym <br> Code | RF- method |
| :---: | :---: | :---: |
| 0.1 | 8.82 | 8.82 |
| 0.2 | 8.30 | 8.28 |
| 0.3 | 7.82 | 7.89 |
| 0.4 | 7.33 | 7.40 |

Table 1 Critical frequencies, in GHz , calculated from the WaveSym code and from the RF method

## Conclusions

A systematic procedure has been given to synthesize consistently with the required operating frequency and injection beam energy, an overmoded sinusoidally corrugated waveguide to be used as a slow-wave structure on backward-wave oscillators. The corrugation parameters are selected so that the waveguide when driven by a $500-\mathrm{keV}$ electron beam operates close to the upper cutoff frequency on the first space harmonic of the lower order TM mode. Consistent with single-mode operation, the periodic structure here designed supports a surface wave that is synchronous with a $500-\mathrm{keV}$ electron beam, and simultaneously exhibits a large coupling impedance, which is a measure of the RF field strength at the beam position. The surface wave is crucial to avoiding mode competition, while large coupling impedance is important for high-efficiency interaction piecewise continuous profile made from a combination of semicircles and rectangles.

## Conclusions

- A sinusoidal profile that meets the design specifications is identified by the corrugation parameters $\mathrm{R}_{0}=4.50 \mathrm{~cm}, \mathrm{~d}=1.4 \mathrm{~cm}$ and $\varepsilon=0.067$ which yield the critical frequency of 8.9 GHz corresponding to the phase velocity of 0.82 c for the synchronous space harmonic.
- The axial electric field strength (and, therefore, the coupling impedance) and the breadth of the passband can be adjusted by proper variation of the corrugation parameters $\varepsilon$ and $\mathrm{R}_{0}$. The deeper the corrugation, the stronger the field and the flatter the passband become.
- The coupling impedance is substancially higher in the piecewise profile than in the sinusoidal counterpart.
- Concerning mode competition, the beam line crosses the dispersion curves of the higher-order modes ( $\mathrm{TM}_{02}$ and $\mathrm{TM}_{03}$ ) at frequencies above 10 GHz , significantly higher than the $8-9 \mathrm{GHz}$ frequency of the nominal $\mathrm{TM}_{01}$ mode.
- Rayleigh-Fourier method has shown excellent performance even when used on the analysis of a piecewise continuous profile made from a combination of semicircles and rectangles

